

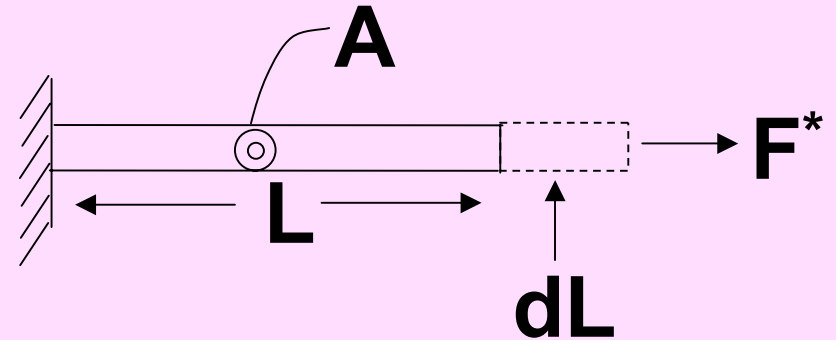
Lecture 3.5

Simple Elastic Systems

SIMPLE ELASTIC SYSTEMS

Basic EQ.

$$Tds \geq dU - f_i'' dx_i'$$



Rev. process

$$f_i'' = f_i' \equiv F^* \quad x_i' \equiv L$$

$$TdS = dU - F^* dL$$

Per unit vol. basis

$$\sigma = \frac{F^*}{A}$$

$$\varepsilon = \frac{L - L_0}{L_0}$$

$$d\varepsilon = \frac{dL}{L_0} - \frac{L}{L_0^2} dL_0 = \frac{dL}{L_0} - (1 + \varepsilon) \frac{dL_0}{L_0}$$

SIMPLE ELASTIC SYSTEMS

Basic Eq. $dU = TdS + F^* dL$

other potentials

$$dH = TdS - L dF^* \quad \{H = U - F^* L$$

$$dF = -SdT + F^* dL \quad \{F = U - TS$$

$$dG = -SdT - L dF^* \quad \{G = H - TS$$

SIMPLE ELASTIC SYSTEMS

Maxwell's Relations

$$\left(\frac{\partial T}{\partial L} \right)_S = \left(\frac{\partial F^*}{\partial S} \right)_\varepsilon$$

$$\left(\frac{\partial T}{\partial F^*} \right)_S = - \left(\frac{\partial L}{\partial S} \right)_{F^*}$$

$$\left(\frac{\partial S}{\partial L} \right)_T = - \left(\frac{\partial F^*}{\partial T} \right)_L$$

$$\left(\frac{\partial S}{\partial F^*} \right)_T = \left(\frac{\partial L}{\partial T} \right)_{F^*}$$

Mnemonic aid :

$$T.S \equiv F^*.L$$

**Sign –ve if T
& F* together**

SIMPLE ELASTIC SYSTEMS

Measurable prop. of Elastic Systems

Sp. heats $C_L = \frac{1}{V_o} \left(\frac{\partial U}{\partial T} \right)_L = \frac{T}{V_o} \left(\frac{\partial S}{\partial T} \right)_L$

$$C_{F^*} = \frac{1}{V_o} \left(\frac{\partial H}{\partial T} \right)_{F^*} = \frac{T}{V_o} \left(\frac{\partial S}{\partial T} \right)_{F^*}$$

Coeff. of Thermal strain $\alpha = \frac{1}{L_o} \left(\frac{\partial L}{\partial T} \right)_{F^*}$

Coeff. of Thermal stress $\beta = \left(\frac{\partial \sigma}{\partial T} \right)_L = \frac{1}{A} \left(\frac{\partial F^*}{\partial T} \right)_L$

Isothermal Young's Modulus $Y = \left(\frac{\partial \sigma}{\partial \varepsilon} \right)_T = \frac{L_o}{A} \left(\frac{\partial F^*}{\partial L} \right)_T$

SIMPLE ELASTIC SYSTEMS

- Interdependence of various prop.

e.g.

$$\alpha = \frac{1}{L_o} \left(\frac{\partial L}{\partial T} \right)_{F^*} = \frac{1}{L_o} \frac{\partial(L, F^*)}{\partial(T, F^*)}$$

$$\beta = \frac{1}{A} \left(\frac{\partial F^*}{\partial T} \right)_L = \frac{\partial(F^*, L)}{\partial(T, L)}$$

$$Y = \frac{L_o}{A} \left(\frac{\partial F^*}{\partial L} \right)_T = \frac{L_o}{A} \frac{\partial(F^*, T)}{\partial(L, T)}$$

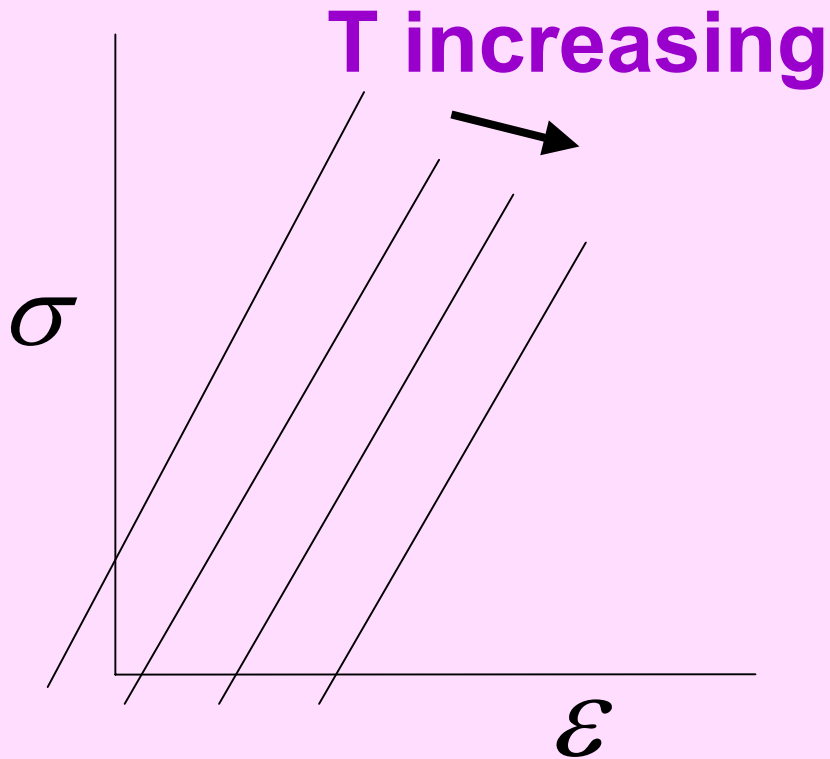
clearly

$$\beta = -Y.\alpha$$

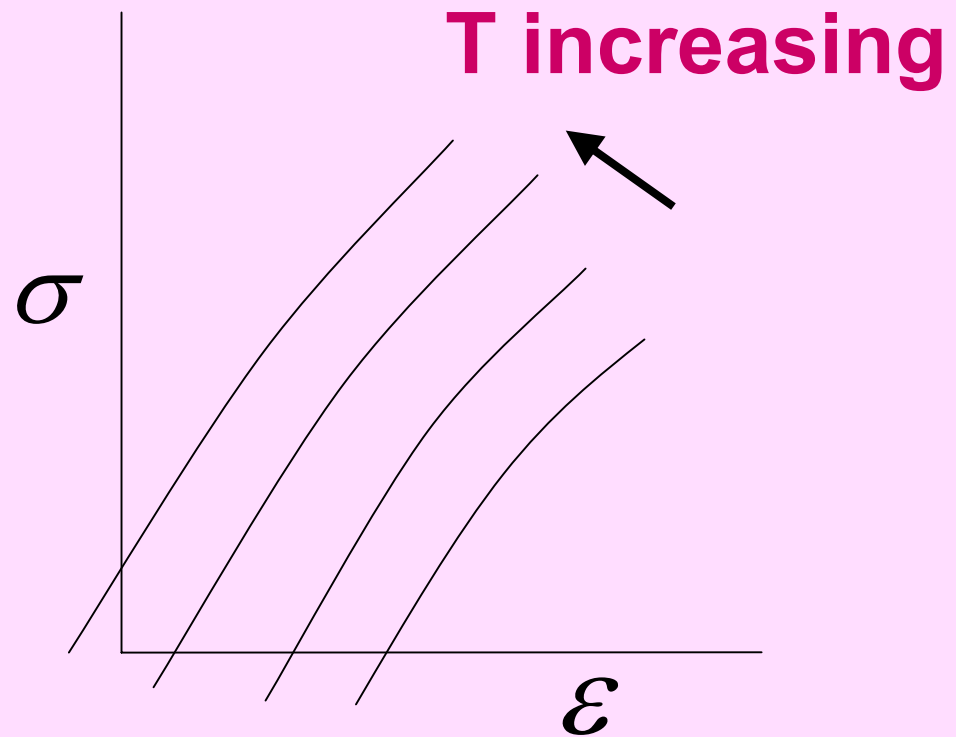
SIMPLE ELASTIC SYSTEMS

Typ. eq. of state

$$\sigma = \sigma(\varepsilon, T)$$



METAL (β -ve)



RUBBER (β +ve)

ANALYSIS OF SIMPLE PROCESSES

Rev. isothermal process

$$S = S(T, F^*)$$

$$dS = \left(\frac{\partial S}{\partial T} \right)_{F^*} dT + \left(\frac{\partial S}{\partial F^*} \right)_T dF^* = \frac{V_o C_{F^*}}{T} dT + \left(\frac{\partial L}{\partial T} \right)_{F^*} dF^*$$

$$= \frac{V_o C_{F^*}}{T} dT + V_o \alpha d\sigma$$

Since the process
is isothermal

$$ds = \alpha d\sigma \text{ where } s = S/V_o$$

ANALYSIS OF SIMPLE PROCESSES

rev isothermal process

\therefore q , Heat Transfer per unit volume is:
{assuming α const.}

$$\int dq = \int T ds = \int T \alpha d\sigma$$
$$= T \cdot \alpha (\sigma_2 - \sigma_1)$$

Work done =

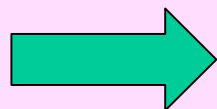
$$\int F^* \cdot dL = \int V_0 \frac{\sigma}{Y} d\sigma$$
$$= \frac{V_0}{2Y} (\sigma_2^2 - \sigma_1^2)$$

{assuming Y const.}

ANALYSIS OF SIMPLE PROCESSES

Rev. adiabatic process

$$dS = 0 = V_o C_{F^*} \frac{dT}{T} + V_o \alpha d\sigma$$


$$\therefore \left(\frac{\partial T}{\partial \sigma} \right)_s = - \frac{T \cdot \alpha}{C_{F^*}}$$

If α , C_{F^*} are assumed constant

$$\Delta T = - \frac{T \alpha}{C_{\sigma}} \cdot \Delta \sigma \quad \text{Assuming } \Delta T \ll T$$

End of Lecture

Lecture 3.6

Simple Elastic Systems.....contd

Solving a few problems

End of Lecture

Lecture 3.7

Simple Magnetic Systems

Simple Magnetic System

Diamagnetism, Para-magnetism, Ferromagnetism

Generalised Force : H^* (magnetic field intensity)

Generalised Coordinate : M (magnetic moment)

M : usually defined / unit volume

Work done $dW = \mu_0 H \cdot V \cdot dM$

$$\mu_0 := 4\pi \times 10^{-7} \text{ N / amp}^2$$

is permeability of free space

Simple Magnetic System

First Law $dQ = dU - \mu_0 V H^* dM$

Second Law $T dS \geq dQ$

Basic Equation $dU = T dS + \mu_0 V H^* dM$

Other potentials $H = U - \mu_0 V H^* M; F = U - TS; G = H - TS$

Simple Magnetic System

Maxwell's Relations

$$dU = TdS + \mu_0 V H^* dM \Rightarrow \left(\frac{\partial T}{\partial M} \right)_S = \mu_0 V \left(\frac{\partial H^*}{\partial S} \right)_M$$

$$dH = TdS - \mu_0 V M dH^* \Rightarrow \left(\frac{\partial T}{\partial H^*} \right)_S = -\mu_0 V \left(\frac{\partial M}{\partial S} \right)_{H^*}$$

$$dF = -SdT + \mu_0 V H^* dM \Rightarrow \left(\frac{\partial S}{\partial M} \right)_T = -\mu_0 V \left(\frac{\partial H^*}{\partial T} \right)_M$$

$$dG = -SdT - \mu_0 V M dH^* \Rightarrow \left(\frac{\partial S}{\partial H^*} \right)_T = \mu_0 V \left(\frac{\partial M}{\partial T} \right)_{H^*}$$

Mnemonic aid ?

Simple Magnetic System

Measurable Properties

Heat Capacities

$$C_{H^*} = \left(\frac{\partial H}{\partial T} \right)_{H^*} = T \left(\frac{\partial S}{\partial T} \right)_{H^*}$$

$$C_M = \left(\frac{\partial U}{\partial T} \right)_M = T \left(\frac{\partial S}{\partial T} \right)_M$$

Simple Magnetic System

Equation of State

Magnetic Susceptibility $\chi = \frac{M}{H^*}$

Simple Equation of State : $\chi V = \frac{C}{T}$

C ∴ Curie's Constant

More accurate : $\chi V = \frac{C}{T - T_0}$

Curie –Weiss equation

(T₀ a constant)

Simple Magnetic System

Reversible Isothermal Process

$$S = S(T, H^*)$$

⇒ **At constant T,**

$$dS = \left(\frac{\partial S}{\partial H^*} \right)_T dH^*$$

Simple Magnetic System

Using Maxwell's relation,

$$dS = \mu_0 V \cdot \left(\frac{\partial M}{\partial T} \right)_{H^*} \cdot dH^*$$

If material obeys

Curie's Equation of State :

$$M = \frac{C}{T} \cdot \frac{H^*}{V} \Rightarrow dS = - \frac{\mu_0 C}{T^2} H^* dH^*$$

SIMPLE MAGNETIC SYSTEMS

Reversible Isothermal.....contd.

Heat Transfer $Q = \int TdS$

$$= -\frac{\mu_0 C}{2T} (H_2^{*2} - H_1^{*2})$$

Thus for $H_2^* > H_1^*$, $Q < 0$ i.e. Heat is rejected during increase in magnetic field

SIMPLE MAGNETIC SYSTEMS

Reversible Adiabatic Process

$$S(T, H^*) \Rightarrow dS = \left(\frac{\partial S}{\partial T} \right)_{H^*} dT + \left(\frac{\partial S}{\partial H^*} \right)_T dH^* = 0$$

Further $\left(\frac{\partial S}{\partial T} \right)_{H^*} = \frac{C_H^*}{T}$; **Substituting this we get**

$$\begin{aligned} \frac{dT}{T} C_{H^*} &= - \left(\frac{\partial S}{\partial H^*} \right)_T dH^* \\ &= -\mu_0 V \left(\frac{\partial M}{\partial T} \right)_{H^*} dH^* \end{aligned}$$

SIMPLE MAGNETIC SYSTEMS

Reversible Adiabatic ProcessContd.

If material obeys Curies' Law

$$M = \frac{C H^*}{T V} \Rightarrow \left(\frac{\partial M}{\partial T} \right)_{H^*} = -\frac{C H^*}{T^2 V}$$

Substituting in earlier equation

$$\frac{dT}{T} C_{H^*} = -\mu_0 V \cdot \left(-\frac{C}{T^2} \cdot \frac{H^*}{V} \right) dH^*$$

$$TdT = \left(\frac{\mu_0 C}{C_H} \right) H^* dH^*$$

SIMPLE MAGNETIC SYSTEMS

Reversible Adiabatic ProcessContd.

Integrating
$$T_2^2 - T_1^2 = \frac{\mu_0 C}{C_{H^*}} (H_2^{*2} - H_1^{*2})$$

Thus, if
$$H_2 < H_1 ; T_2 < T_1$$

i.e. we can reduce temperature by adiabatic demagnetisation right upto Curie Point.

Used in reaching very low temperatures

Solving a few problems

INTRODUCTION TO COMPLEX SYSTEMS

COMPLEX SYSTEMS

⇒ More than one mode of work
interaction

**Heat Transfer = Change in Int. Energy
– Work input**

$$TdS = dU - f_1 dx_1 - f_2 dx_2$$

COMPLEX SYSTEMS

Examples

$$TdS = dU + PdV - V_e dQ_e$$

**Reversible Cell
Liberating Gases**

$$TdS = dU + PdV - \mu_0 V H^* dM$$

**Magnetization of
Compressible
Substances**

$$TdS = dU - F^* dL - E d\Pi$$

**Piezo-electric
System**