

# Molecular Optics and Thin Films



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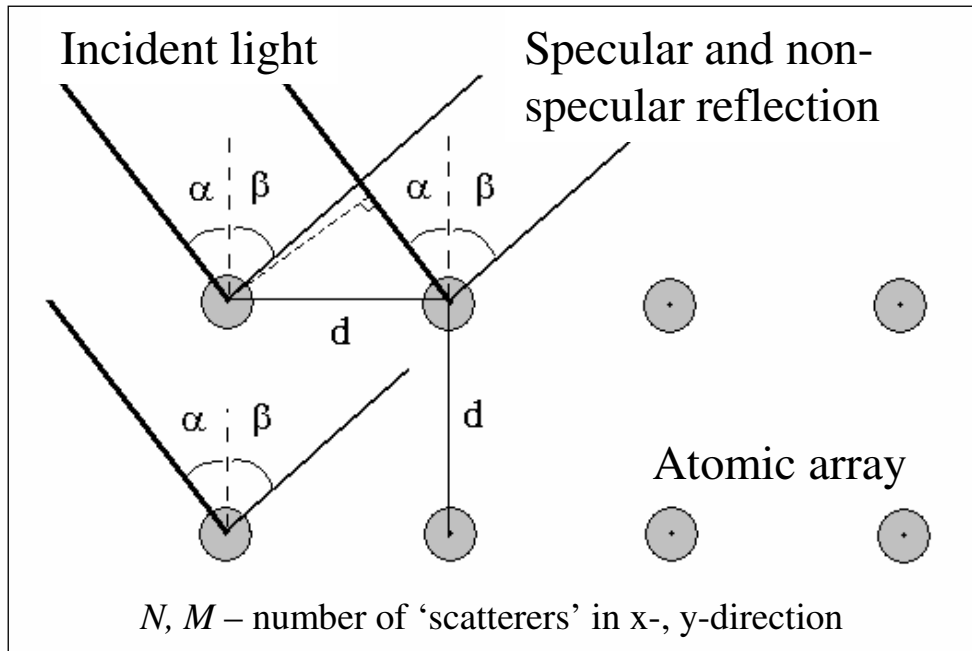


Our studies address microscopic calculations of the optical properties of thin films in the soft x-ray and extreme ultraviolet (EUV) regime when the wavelength is on order of the inter-atomic distances. Investigations include the effects of surface roughness, as well as one-dimensional “amorphous” atomic spacing, on specular and non-specular reflection.

# When wavelength gets small...

- When approaching reflection and transmission macroscopically – i.e. when the wavelength of light  $\lambda \gg d$ , the inter-atomic distance – determining optical properties becomes an optics problem. The interface is, for all practical purposes, “flat.”
- When  $\lambda$  is on the same order as  $d$  (the EUV and x-ray regime), the interface no longer “looks” flat, but is more correctly treated as an array of scattering atoms, microscopically speaking.
- Our task: Investigate how microscopic atomic properties such as position correlations and scattering amplitudes affect macroscopic reflection
  - The field of Molecular Optics!
  - Side note: BYU’s Thin Films research group is experimentally determining optical properties in the EUV regime for  $\text{UO}_x$ . Could the theoretical microscopic scattering amplitudes be correlated to these experimental results?

# Diagrams and Definitions



The symmetry between the x and y directions has been verified via theory and program, as was expected

Path length difference between two ‘atoms’

$$\delta = kd(\sin \alpha - \sin \beta)$$

Expression in 1D, assuming isotropic [s-wave] scattering (including constant scattering probability factor  $f$ ):

$$Field = \sum_{n=0}^N f e^{in\delta}$$

$$I(\delta, N, f) = \frac{|f|^2 \cos(\delta N + \delta) - 1}{N \cos(\delta) - 1}$$

Extended to 2D:

$$Field = \sum_{m=0}^{M-1} \left( \frac{e^{im\Delta} f e^{iN\delta}}{e^{i\delta} - 1} - \frac{f e^{im\Delta}}{e^{i\delta} - 1} \right) = f^2 \left( \frac{e^{iN\delta} - 1}{e^{i\delta} - 1} \right) \left( \frac{e^{iM\Delta} - 1}{e^{i\Delta} - 1} \right)$$

looking strangely enough like a separable equation, where

$$\Delta = kd(\cos \alpha + \cos \beta)$$

# Dependent Variables

- The path length differences  $\delta$  and  $\Delta$  are dependent on atomic spacing  $d$ , wave number  $k$ , incident angle  $\alpha$ , and exit angle  $\beta$
- That is to say,

$$\delta(d,k,\alpha,\beta) \text{ and } \Delta(d,k,\alpha,\beta),$$

$$E(N,M,d,k,\alpha,\beta) \text{ and } I(N,M,d,k,\alpha,\beta),$$

where  $I = |E|^2$

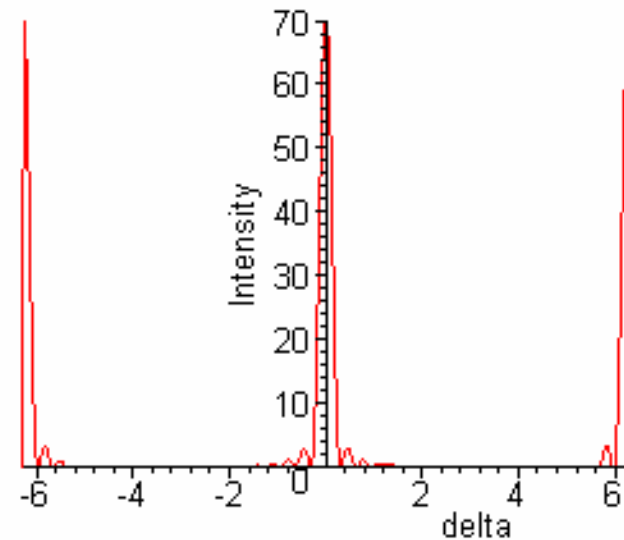
- To more correctly model  $\Delta$ , an attenuation factor accounting for scattering from each succeeding surface,  $\xi$ , was introduced
- For simplicity, lumped into  $\Delta$ , where  $\Delta \rightarrow \Delta + i\xi$ 
  - As expected, the intensity was affected

# Thickness Effects

For two-dimensions:

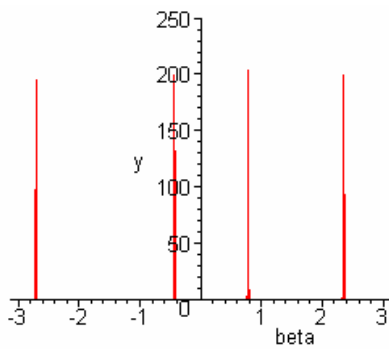
What happens to  $I(\delta)$  as  $M$  increases?

- Intensity increases to a maximum, then decreases to zero, then increases to a lesser maximum, exhibiting features of a dying sinusoidal exponential
- The first intensity max and min occurs at  $M=6$  and  $16$ , respectively.

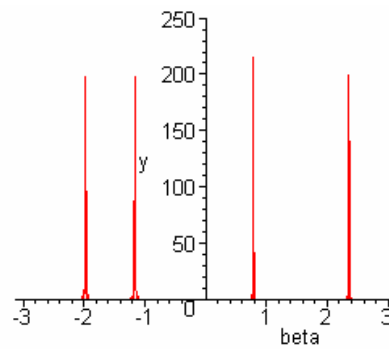


Parameter	Value	Rows ( $M$ )	Intensity
$N$	100	1	20
$d$ in x-dir	135	6	75 (max)
$d$ in y-dir	$d_x$	16	0
$\Delta$	$\pi/8$	23	22
-	-	32	0
-	-	40	13
-	-	48	0

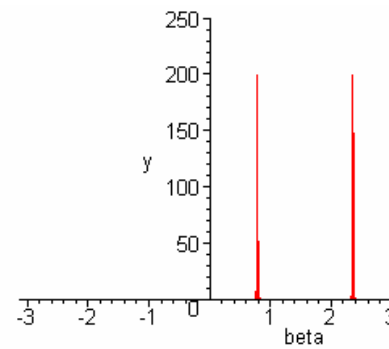
# Wavelength Scan



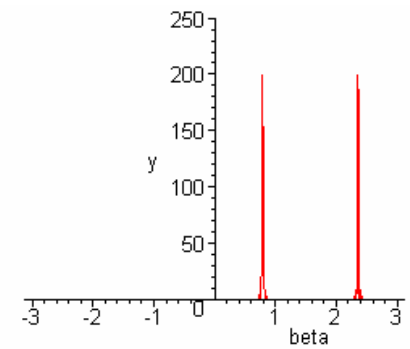
$$\lambda = 11d/8$$



$$\lambda = 13d/8$$



$$\lambda = 14d/8$$

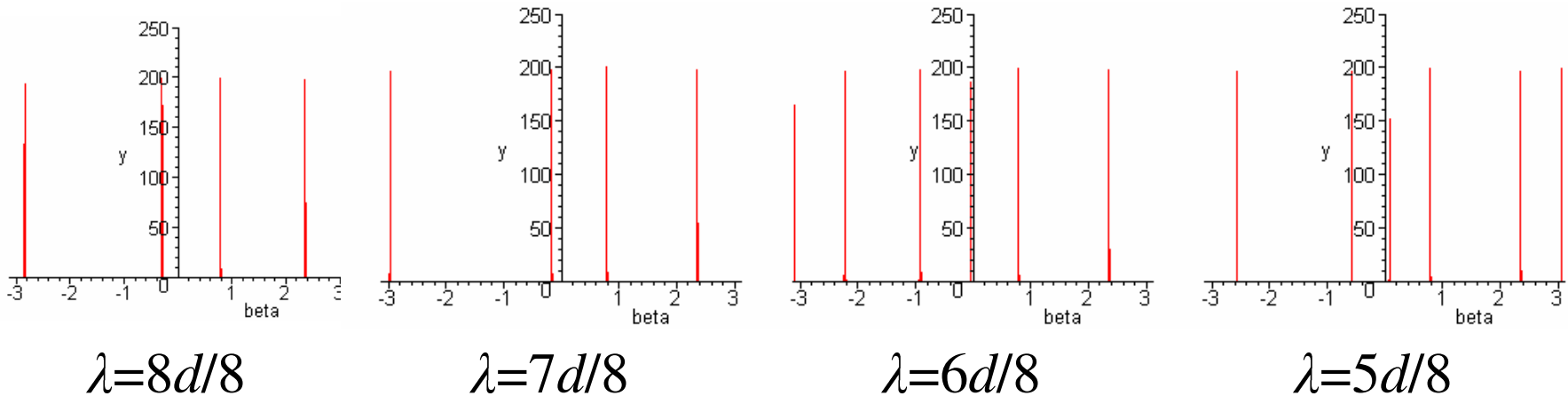


$$\lambda = 30d/8$$

In each case,  
 $N=200$   
 $M=1$   
 $d=135$   
 $\alpha=\pi/4$

- As  $\lambda$  increases with respect to  $d$ ,  $I(\beta)$  from  $-\pi < \beta < \pi$  shifts toward  $\sim 1.7$  and then diminishes
- There comes a point where, for  $\alpha = \pi/4$ ,  $\beta$  is only  $\pi/4$  and  $3\pi/4$ , where either the incident light follows Snell's Law or passes right through the lattice.

# Wavelength Scan

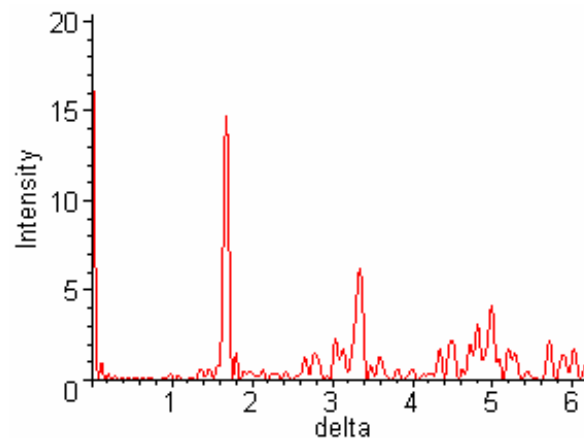
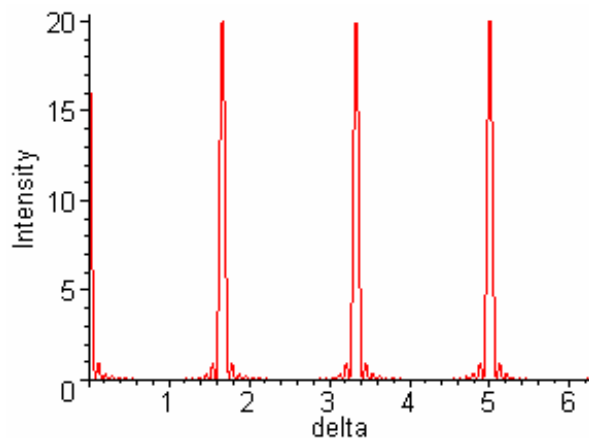


In each case,  
 $N=200$   
 $M=1$   
 $d=135$   
 $\alpha=\pi/4$

- Intensity as a function of  $\beta$
- Out of one intensity peak, many, and vice versa
  - As  $\lambda$  increases, the specular spikes move, multiply, and/or decrease

# Adding Randomness

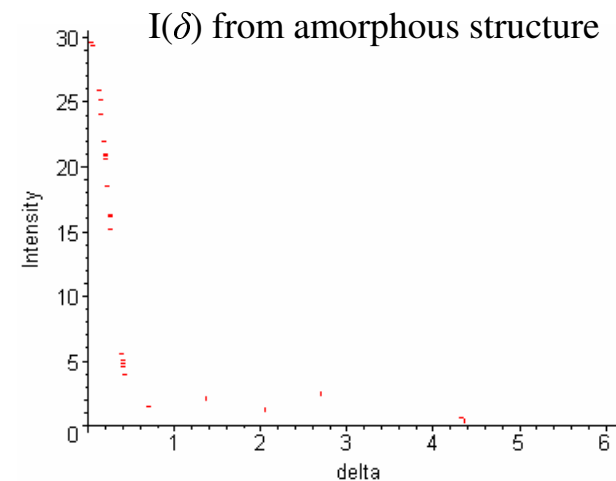
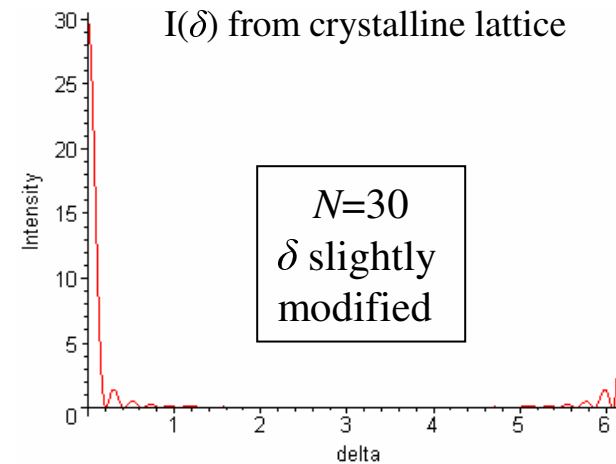
- To better understand the thin film aspect, a Gaussian random number generator was added to take the inter-atomic spacing from the crystalline case to the amorphous case
  - Sputtered thin films are generally amorphous in nature
- Atoms were shifted from their *usual places in the lattice*
- Note: The results are in terms of path length difference  $\delta$



$N=10$   
 $M=1$   
 $d=120$   
 $d+(-1,0,1)$   
 $\lambda=100$

# Amorphous Atomic Arrangement (AAA)

- Variations in atomic spacing  $d$  simulated by a random number generator
  - Atomic shift for each atom measured *relative to the previous neighboring atom's shift*
- To what extent is diffraction preserved?
- The amorphous plot required more generating of points
  - Why this is yet unknown



# Conclusions

- Most of our results are still preliminary, but with continued research we hope to answer the question of theoretical optical constant determination. We have:
  - reproduced specular and Bragg reflections from isotropic scattering processes
  - showed the effects of disorder on the width and height of the Bragg peaks
- We may be able to:
  - predict the order in the atomic positions through height comparison of specular and Bragg peaks
  - perform qualitative studies of the effects of different kinds of roughness on the reflection coefficient of the surface
  - Perform a qualitative study of different scattering cross-sections on reflection
- It may be possible to theoretically calculate *accurate* optical constants for elements in the EUV and soft x-ray regime
  - This refining of optical constants could have major impacts in industries such thin films, affecting everything from EUV lithography to satellite-based observatories

# References

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Notes from interviews with Dr. Turley. Late March to early September, 2002.