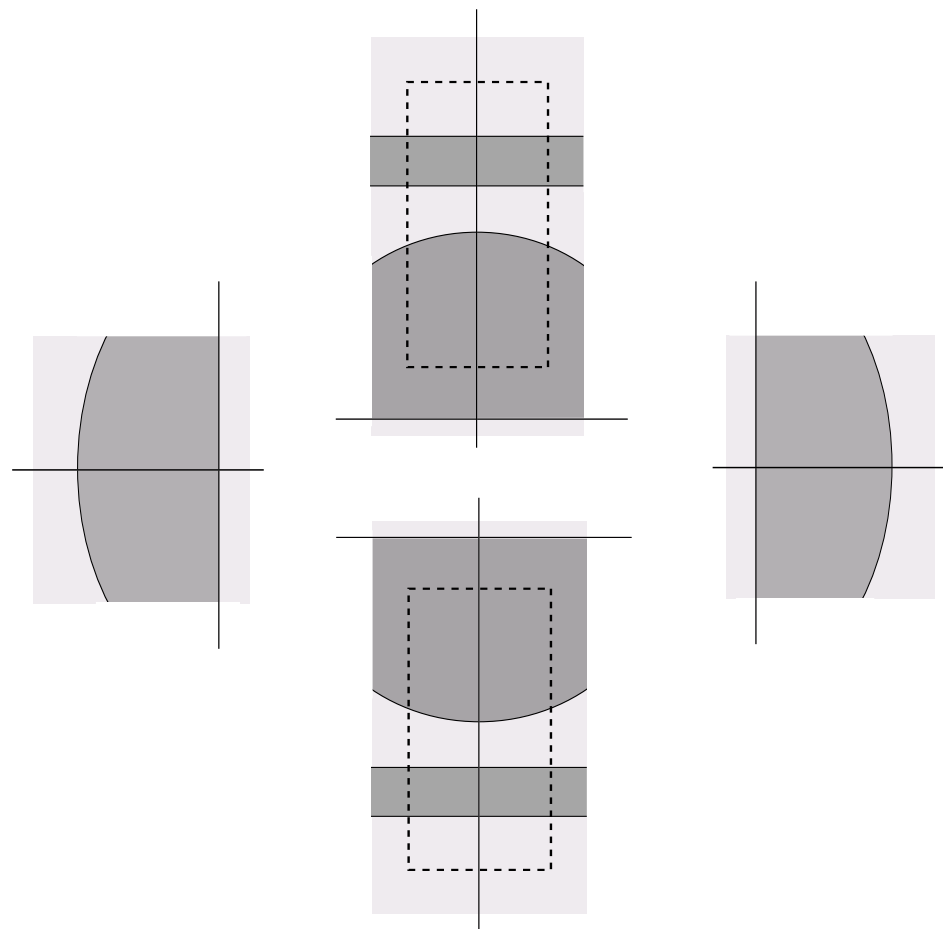
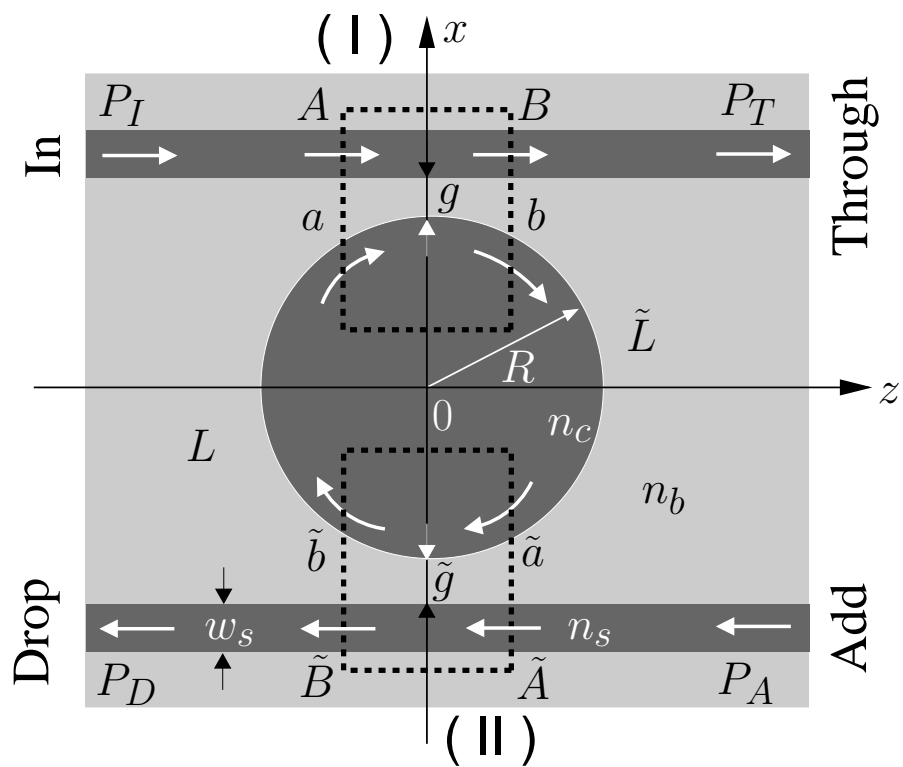


Perturbational evaluation of bend mode phase shifts for the tuning of cylindrical microresonators

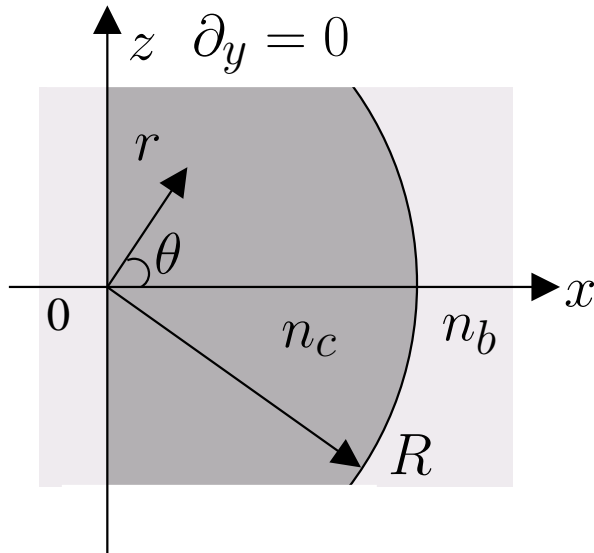


Kirankumar Hiremath, Manfred Hammer
Applied Analysis and Mathematical Physics Group
Department of Applied Mathematics
University of Twente, The Netherlands

Schematic representation of 2D microresonator



Bent Waveguide



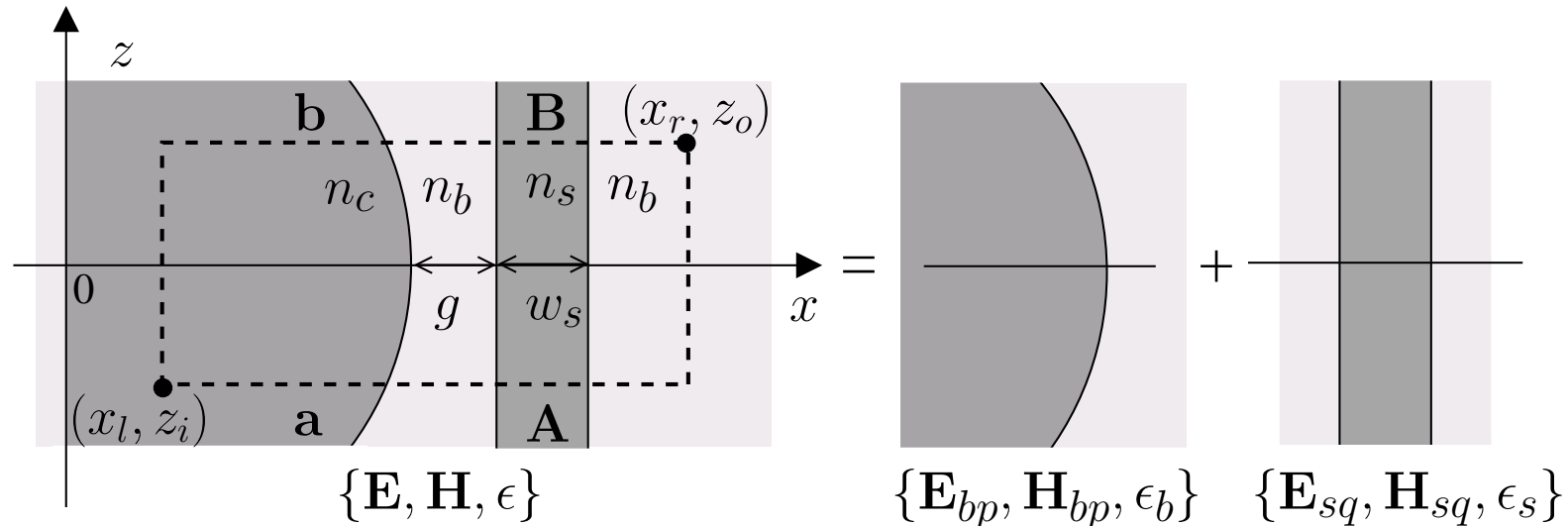
- Field ansatz:

$$\mathbf{E} = (\tilde{E}_r(r), \tilde{E}_y(r), \tilde{E}_\theta(r)) e^{i(\omega t - \gamma R \theta)}$$

$$\mathbf{H} = (\tilde{H}_r(r), \tilde{H}_y(r), \tilde{H}_\theta(r)) e^{i(\omega t - \gamma R \theta)}$$

- $\omega \in \mathbb{R}$
- Propagation constant $\gamma = \beta - i\alpha \in \mathbb{C}$
- Basic components \tilde{E}_y (TE) and \tilde{H}_y (TM)
- *Bessel equation* for \tilde{E}_y and \tilde{H}_y of order γR
- Piecewise field ansatz + Interface conditions \Rightarrow Dispersion equation
- Dispersion equation + Complex order Bessel functions \Rightarrow **Analytical Bend Modes**

Bent-straight waveguide coupler

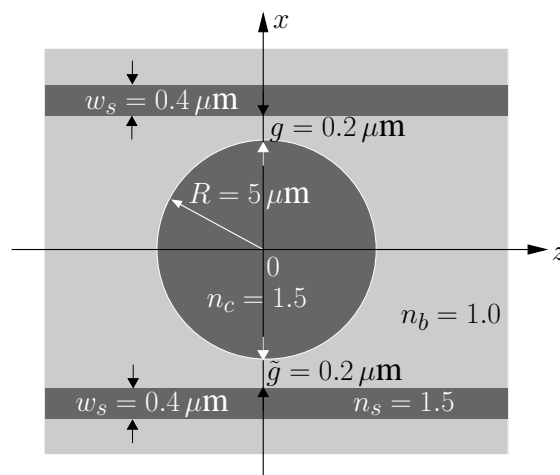
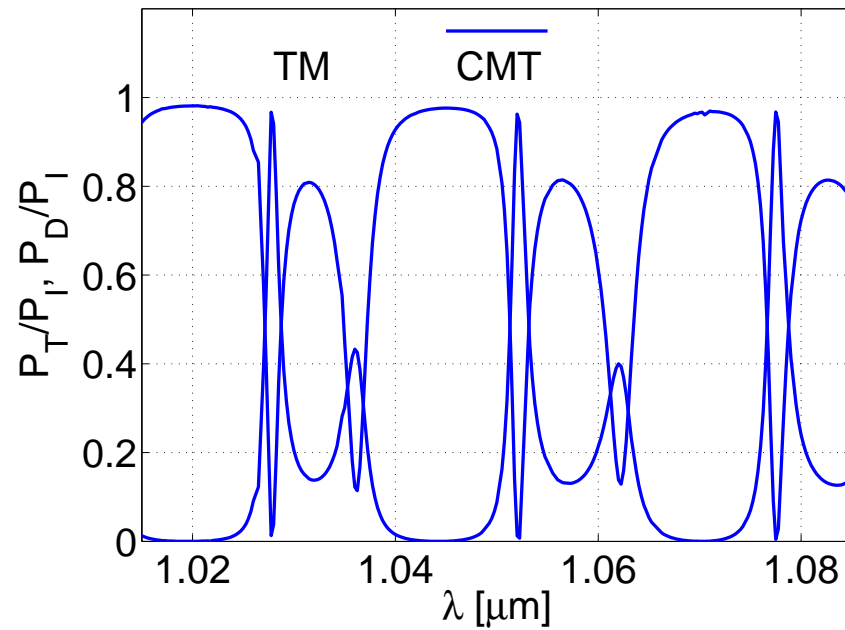
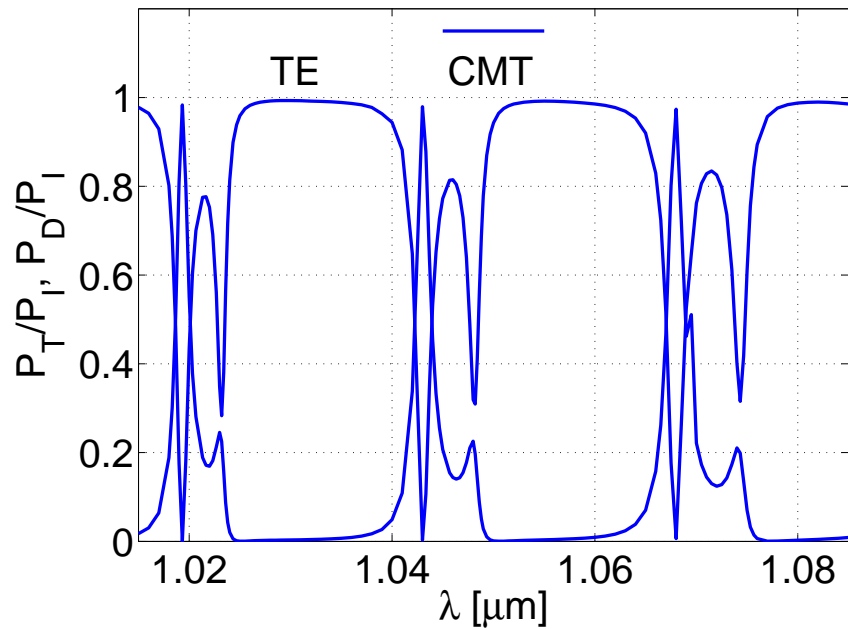


- Coupled Mode ansatz: C_{bp}, C_{sq} unknown amplitudes

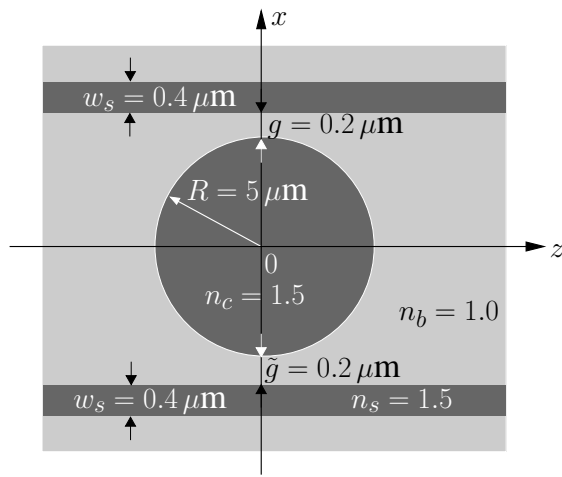
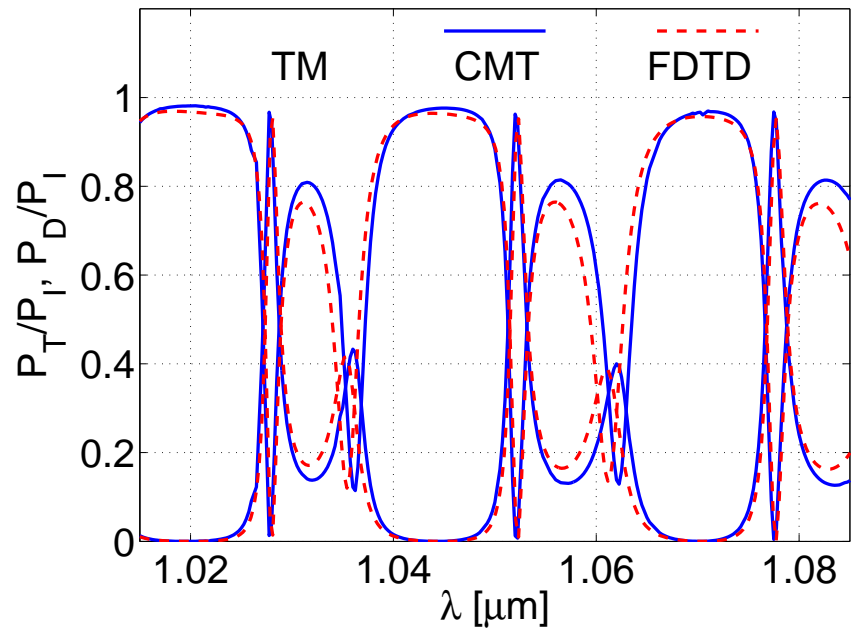
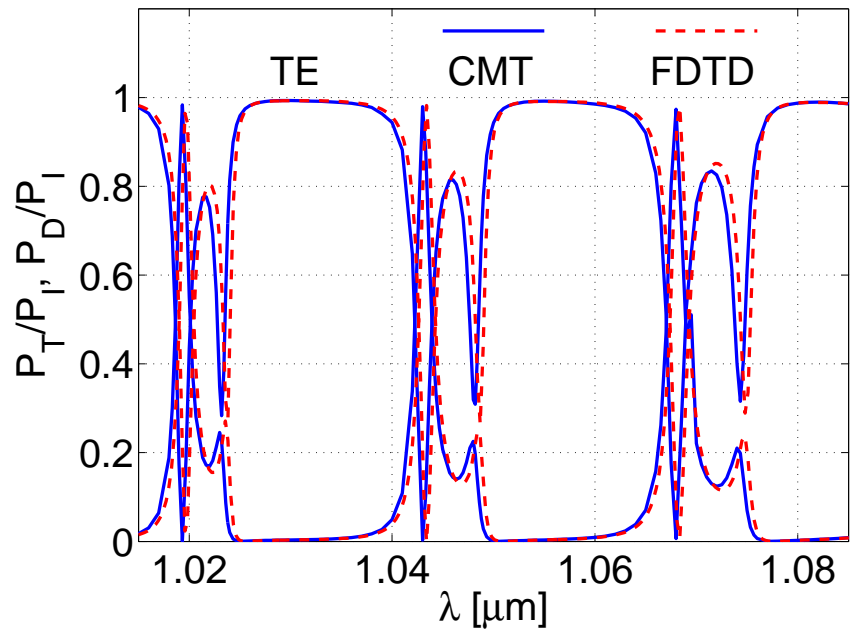
$$\begin{pmatrix} \mathbf{E}(x, z) \\ \mathbf{H}(x, z) \end{pmatrix} = \sum_{p=1}^{N_b} C_{bp}(z) \begin{pmatrix} \mathbf{E}_{bp}(x, z) \\ \mathbf{H}_{bp}(x, z) \end{pmatrix} + \sum_{q=1}^{N_s} C_{sq}(z) \begin{pmatrix} \mathbf{E}_{sq}(x, z) \\ \mathbf{H}_{sq}(x, z) \end{pmatrix}$$

- Lorentz reciprocity theorem / variational method $\mapsto \mathbf{M}(z) \cdot d_z \mathbf{C}(z) = \mathbf{F}(z) \cdot \mathbf{C}(z)$
- Solve for \mathbf{C} + projection on straight waveguide modes \mapsto scattering matrix (\mathbf{S})

Spectral response

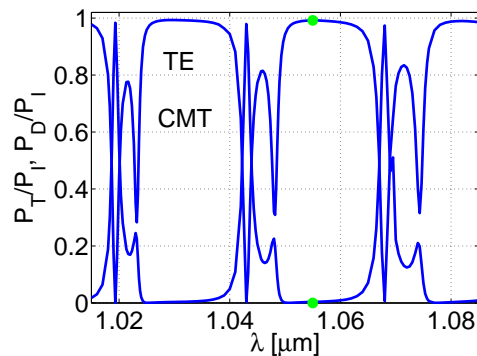
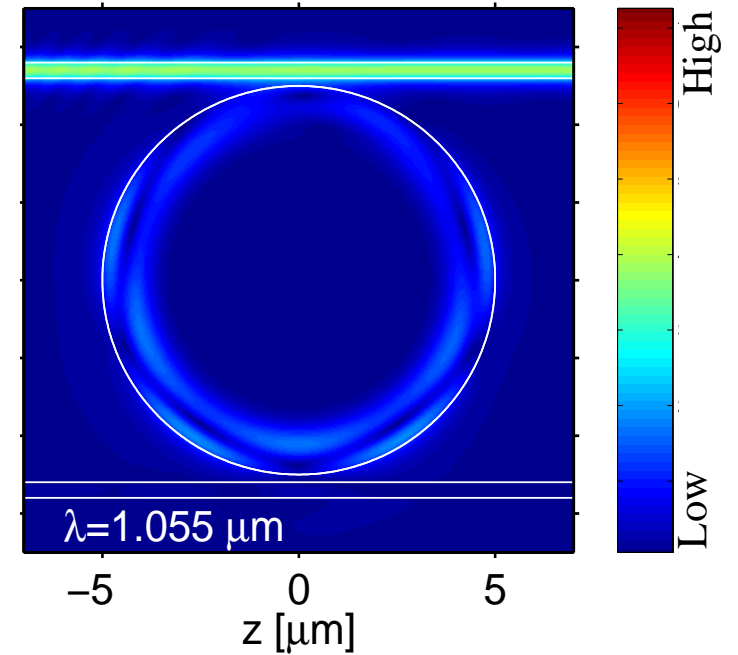


Spectral response



Microresonator field profiles

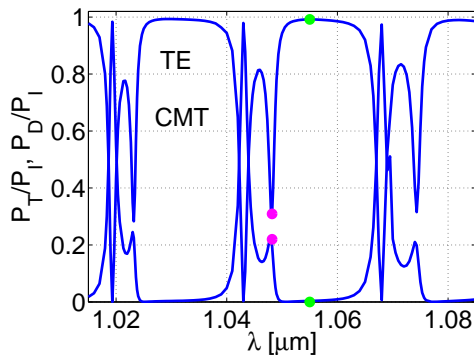
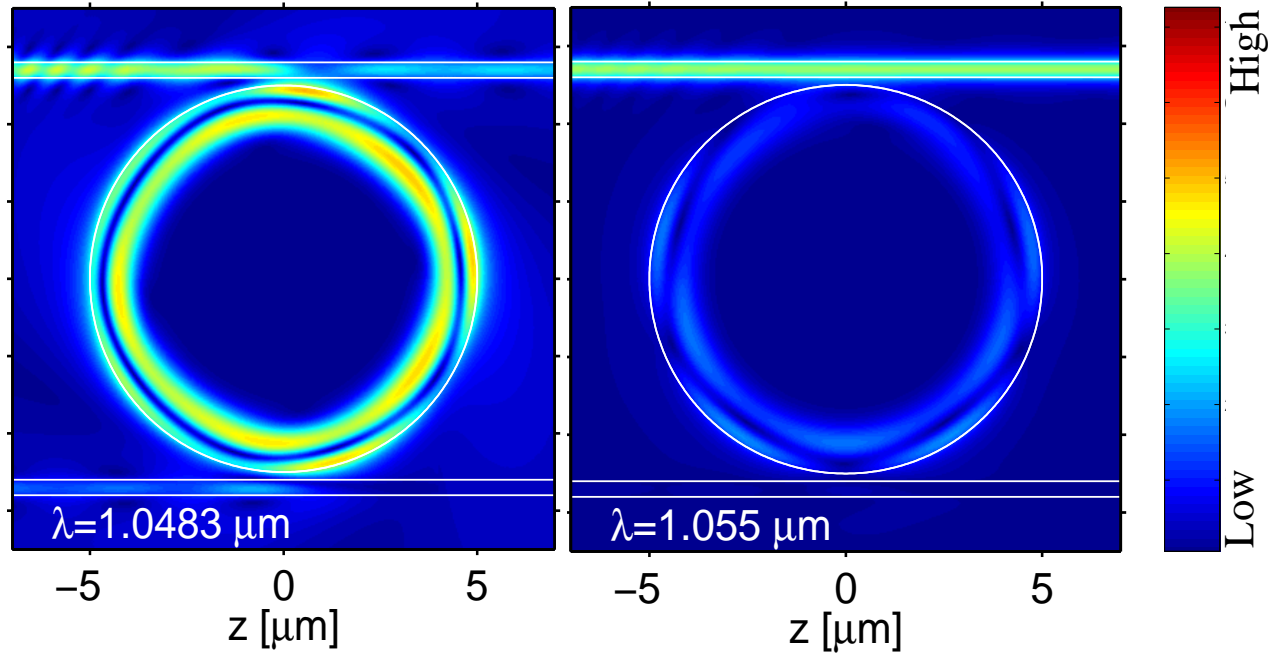
- TE $|E_y|$: Coupled mode theory microresonator simulations with 3 cavity modes



- $\lambda = 1.055 \mu\text{m} \mapsto$ off resonance

Microresonator field profiles

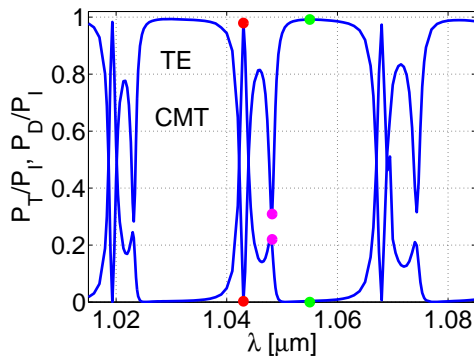
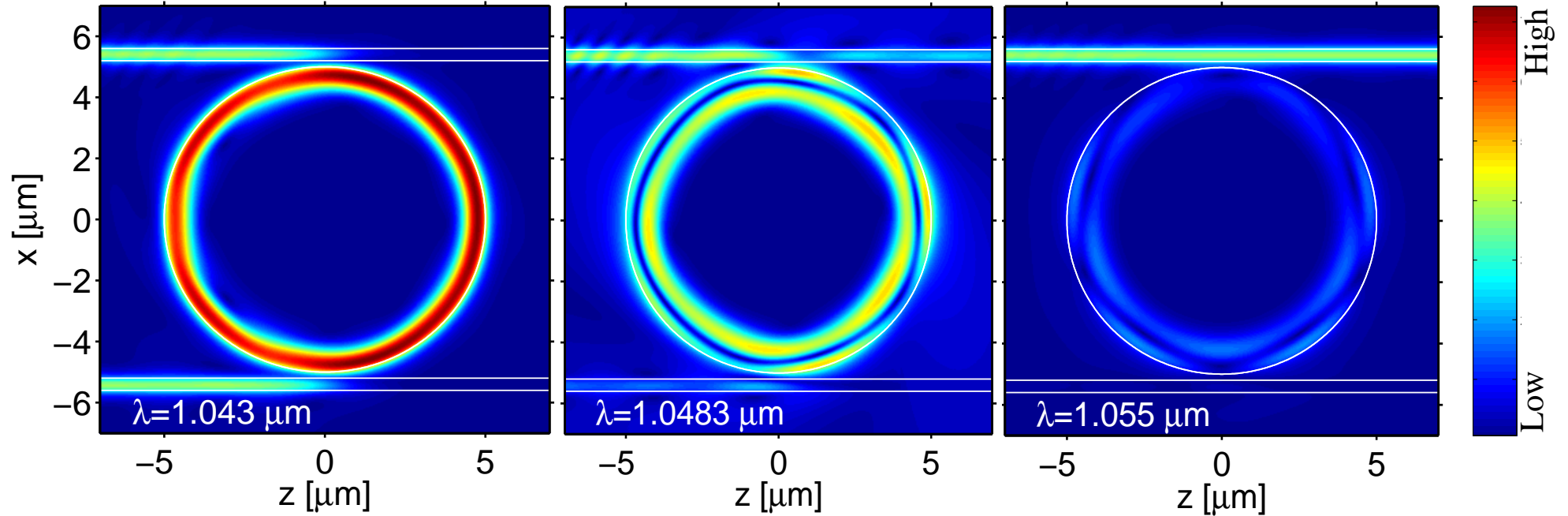
- TE $|E_y|$: Coupled mode theory microresonator simulations with 3 cavity modes



- $\lambda = 1.0483 \mu\text{m} \mapsto$ Resonance corresponding to first order mode
- $\lambda = 1.055 \mu\text{m} \mapsto$ off resonance

Microresonator field profiles

- TE $|E_y|$: Coupled mode theory microresonator simulations with 3 cavity modes



- $\lambda = 1.043 \mu\text{m} \mapsto$ Resonance corresponding to fundamental mode
- $\lambda = 1.0483 \mu\text{m} \mapsto$ Resonance corresponding to first order mode
- $\lambda = 1.055 \mu\text{m} \mapsto$ off resonance

Tuning

- Over a small wavelength interval, the scattering matrix is assumed to be constant. **Phase relations** of the cavity modes \leftrightarrow **resonances**
- At resonance $2\pi R\beta_m + \phi = 2\pi m$
 ϕ : Phase change due to straight waveguide coupling
- Tuning: p perturbation (tuning) parameter

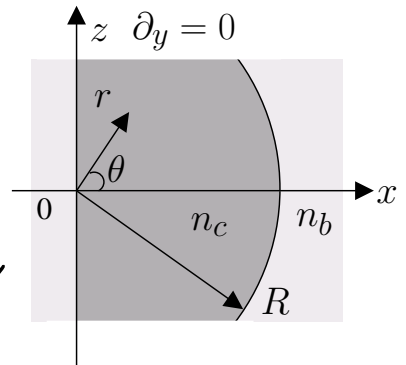
$$\beta(p, \tilde{\lambda}_m) \approx \beta(0, \lambda_m) + p \left. \frac{\partial \beta}{\partial p} \right|_{0, \lambda_m} + (\tilde{\lambda}_m - \lambda_m) \left. \frac{\partial \beta}{\partial \lambda} \right|_{0, \lambda_m} \stackrel{!}{=} \beta_m$$

$$\text{Shift due to tuning : } \Delta\lambda = p \left. \left(\frac{\partial \beta}{\partial p} \right) \left(\frac{\partial \beta}{\partial \lambda} \right)^{-1} \right|_{0, \lambda_m}$$

- Homogeneity arguments: $\frac{\partial \beta}{\partial \lambda} \approx -\frac{\beta}{\lambda}$
- For cavity core permittivity perturbation $p = \delta\epsilon_c \quad \mapsto \quad$ **How to find $\frac{\partial \beta}{\partial \epsilon_c}$?**

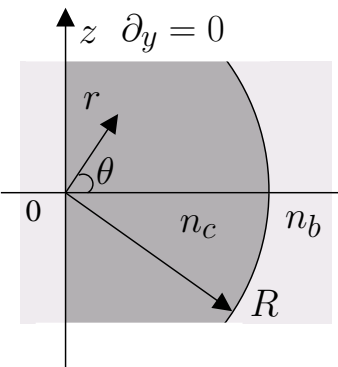
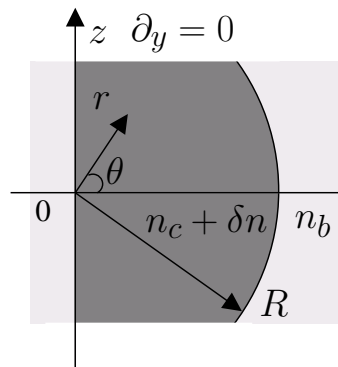
Permittivity perturbation

$$\underbrace{\begin{pmatrix} \mathbf{E} \\ \mathbf{H} \end{pmatrix}}_{\text{Unperturbed } n_c = \sqrt{\epsilon_c}} = \begin{pmatrix} \tilde{\mathbf{E}} \\ \tilde{\mathbf{H}} \end{pmatrix} (r) e^{-i\gamma R\theta}$$



Permittivity perturbation

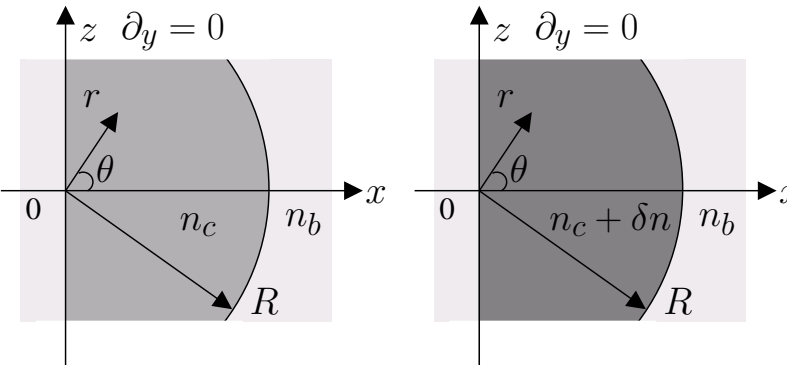
$$\underbrace{\begin{pmatrix} \mathbf{E} \\ \mathbf{H} \end{pmatrix}}_{\text{Unperturbed } n_c = \sqrt{\epsilon_c}} = \begin{pmatrix} \tilde{\mathbf{E}} \\ \tilde{\mathbf{H}} \end{pmatrix} (r) e^{-i\gamma R\theta}$$

$$\underbrace{\begin{pmatrix} \mathbf{E}_p \\ \mathbf{H}_p \end{pmatrix}}_{\text{Perturbed } n_c + \delta n = \sqrt{\epsilon_p}} = C(\theta) \begin{pmatrix} \mathbf{E} \\ \mathbf{H} \end{pmatrix}$$

14

Permittivity perturbation

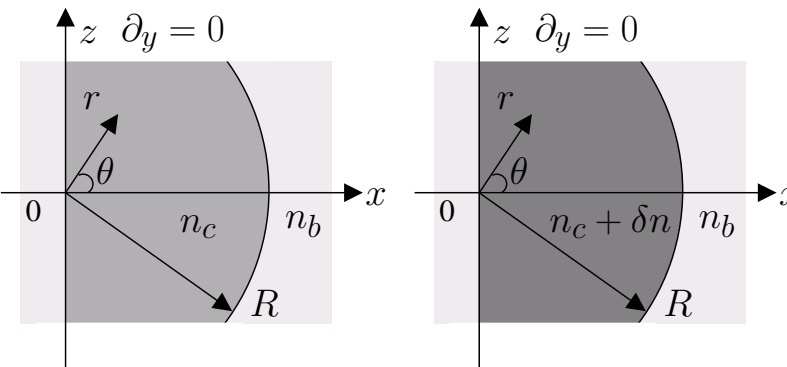
$$\underbrace{\begin{pmatrix} \mathbf{E} \\ \mathbf{H} \end{pmatrix}}_{\text{Unperturbed } n_c = \sqrt{\epsilon_c}} = \begin{pmatrix} \tilde{\mathbf{E}} \\ \tilde{\mathbf{H}} \end{pmatrix} (r) e^{-i\gamma R\theta}$$


$$\underbrace{\begin{pmatrix} \mathbf{E}_p \\ \mathbf{H}_p \end{pmatrix}}_{\text{Perturbed } n_c + \delta n = \sqrt{\epsilon_p}} = C(\theta) \begin{pmatrix} \mathbf{E} \\ \mathbf{H} \end{pmatrix}$$

- **Lorentz's reciprocity theorem:** $(\mathbf{E}_p, \mathbf{H}_p, \epsilon_p)$ and $(\mathbf{E}, \mathbf{H}, \epsilon_c) \mapsto \mathbf{M} \cdot d_\theta C(\theta) = \mathbf{F} \cdot C(\theta)$

$$\begin{pmatrix} \mathbf{E}_p \\ \mathbf{H}_p \end{pmatrix} = C_0 \begin{pmatrix} \tilde{\mathbf{E}} \\ \tilde{\mathbf{H}} \end{pmatrix} (r) e^{-i(\gamma + \delta\gamma)R\theta} \qquad \delta\gamma = \frac{\omega\epsilon_0}{R} \frac{\int_0^R (\epsilon_p - \epsilon_c) \mathbf{E} \cdot \mathbf{E}^* r dr}{\int_0^\infty \mathbf{a}_\theta \cdot (\mathbf{E} \times \mathbf{H}^* + \mathbf{E}^* \times \mathbf{H}) dr}$$

Permittivity perturbation

$$\underbrace{\begin{pmatrix} \mathbf{E} \\ \mathbf{H} \end{pmatrix}}_{\text{Unperturbed } n_c = \sqrt{\epsilon_c}} = \begin{pmatrix} \tilde{\mathbf{E}} \\ \tilde{\mathbf{H}} \end{pmatrix} (r) e^{-i\gamma R\theta}$$


$$\underbrace{\begin{pmatrix} \mathbf{E}_p \\ \mathbf{H}_p \end{pmatrix}}_{\text{Perturbed } n_c + \delta n = \sqrt{\epsilon_p}} = C(\theta) \begin{pmatrix} \mathbf{E} \\ \mathbf{H} \end{pmatrix}$$

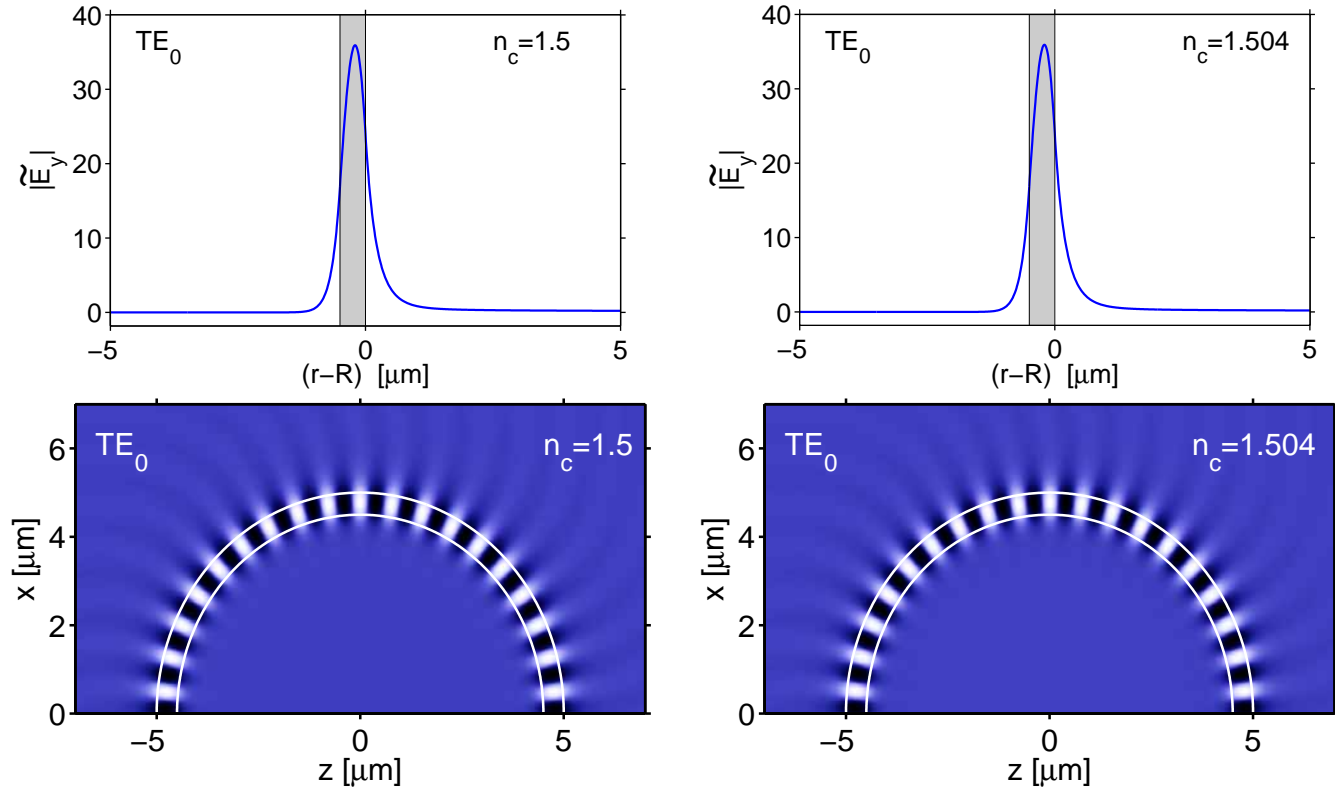
- Lorentz's reciprocity theorem:** $(\mathbf{E}_p, \mathbf{H}_p, \epsilon_p)$ and $(\mathbf{E}, \mathbf{H}, \epsilon_c) \mapsto \mathbf{M} \cdot d_\theta C(\theta) = \mathbf{F} \cdot C(\theta)$

$$\begin{pmatrix} \mathbf{E}_p \\ \mathbf{H}_p \end{pmatrix} = C_0 \begin{pmatrix} \tilde{\mathbf{E}} \\ \tilde{\mathbf{H}} \end{pmatrix} (r) e^{-i(\gamma + \delta\gamma)R\theta} \quad \delta\gamma = \frac{\omega\epsilon_0}{R} \frac{\int_0^R (\epsilon_p - \epsilon_c) \mathbf{E} \cdot \mathbf{E}^* r dr}{\int_0^\infty \mathbf{a}_\theta \cdot (\mathbf{E} \times \mathbf{H}^* + \mathbf{E}^* \times \mathbf{H}) dr}$$

- $\delta\gamma$ is real \implies change in phase constant ($\delta\beta$)

$$\frac{d\beta}{d\epsilon_c} \approx \frac{\delta\beta}{\delta\epsilon_c} = \frac{\omega\epsilon_0}{R} \frac{\int_0^R \tilde{\mathbf{E}} \cdot \tilde{\mathbf{E}}^* r dr}{\int_0^\infty \mathbf{a}_\theta \cdot (\tilde{\mathbf{E}} \times \tilde{\mathbf{H}}^* + \tilde{\mathbf{E}}^* \times \tilde{\mathbf{H}}) dr}$$

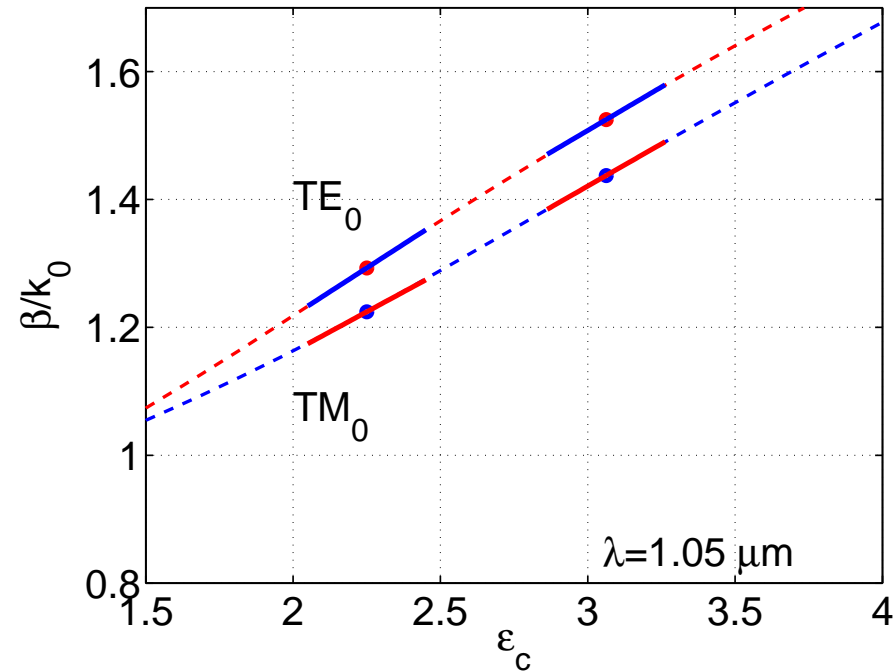
Bend modes



$|\tilde{E}_y|$ and real physical E_y for a bent waveguide $n_b = 1.0$, width = $0.5 \mu\text{m}$, $R = 5 \mu\text{m}$, $\lambda = 1.05 \mu\text{m}$.

	$n_c = 1.5$	$n_c = 1.504$	
	N_{eff} (Direct)	N_{eff} (Direct)	$\Re(N_{\text{eff}})$ (Perturb.)
TE ₀	$1.29297 - i 7.5205 \cdot 10^{-6}$	$1.29655 - i 6.4158 \cdot 10^{-6}$	1.2966

Bend mode phase shifts - I

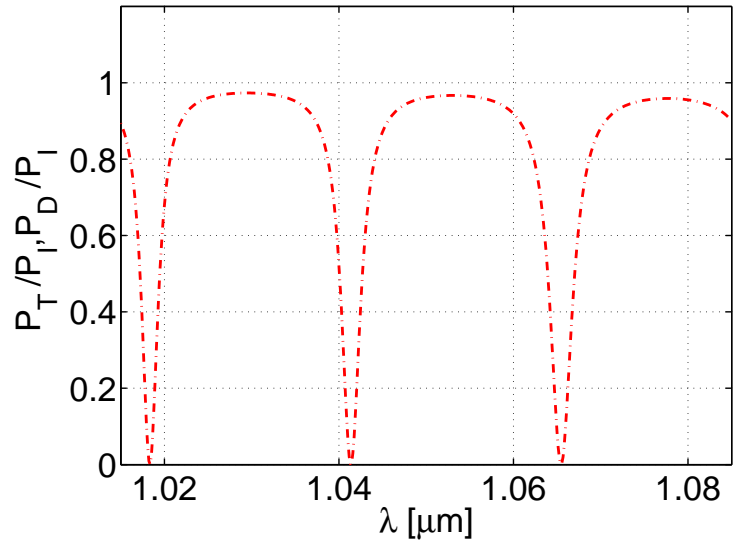


Dashed lines = β/k_0 direct calculations

Dots = Reference points ($\epsilon_c = (1.5)^2$ and $\epsilon_c = (1.75)^2$)

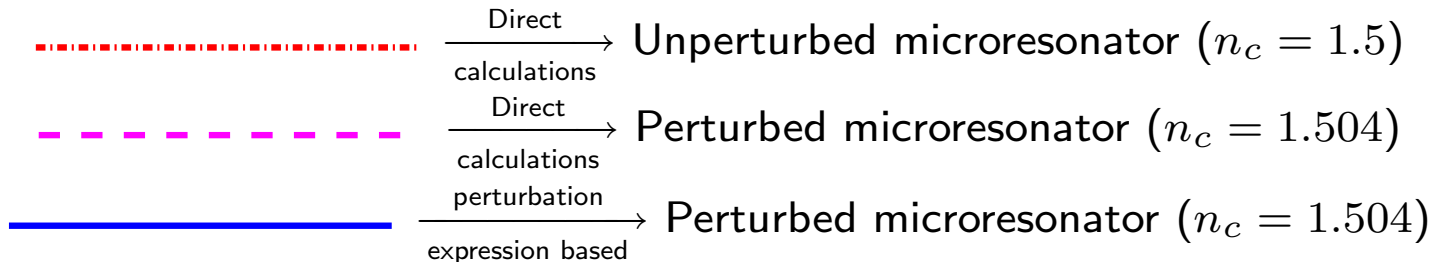
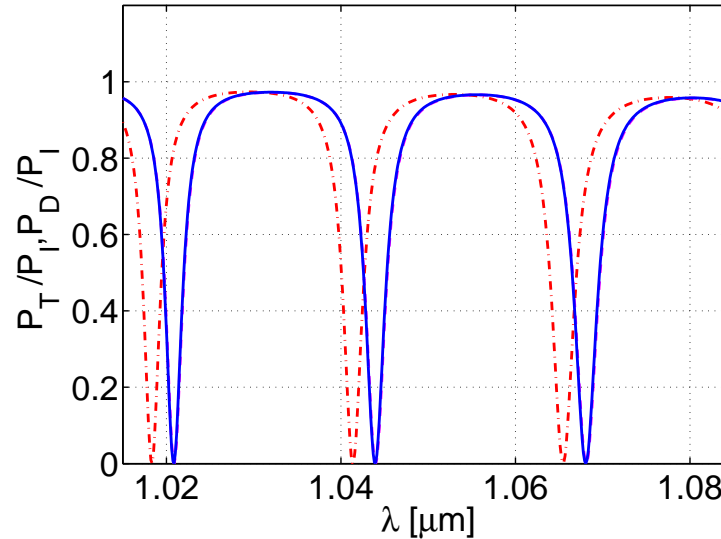
Solid segments = $\beta_{ref} + \delta\epsilon_c \left. \frac{\partial\beta}{\partial\epsilon_c} \right|_{ref}$, perturbational evolution of $\frac{\partial\beta}{\partial\epsilon_c}$

Microresonator spectrum perturbational evaluation - I



----- Direct
calculations → Unperturbed microresonator ($n_c = 1.5$)

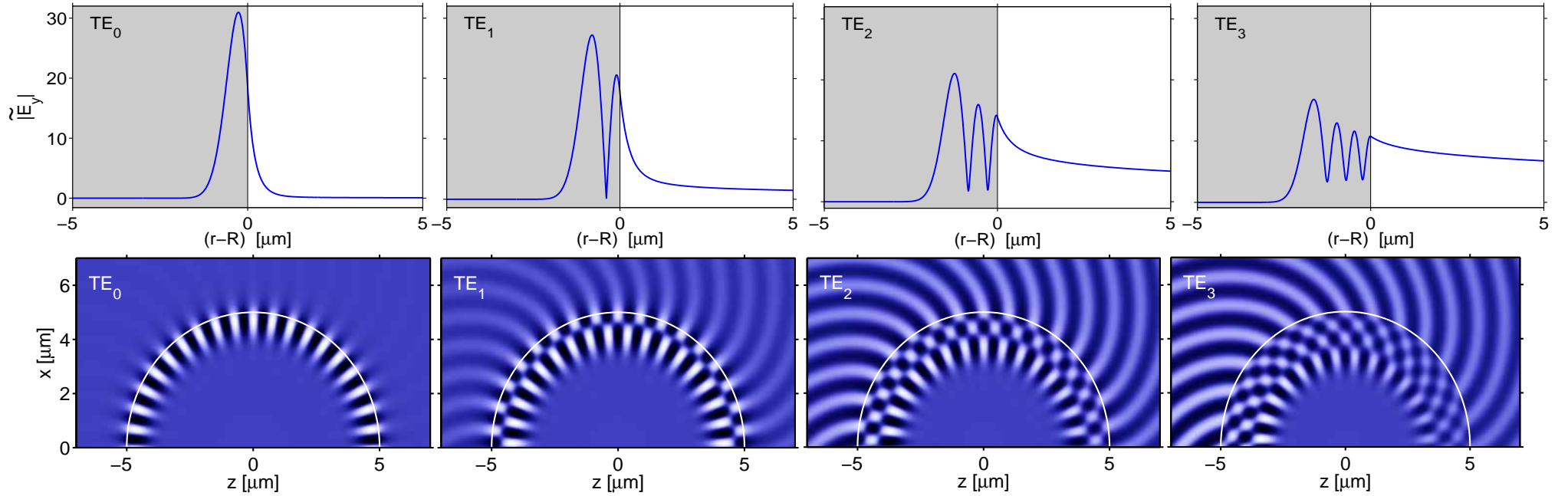
Microresonator spectrum perturbational evaluation - I



- Resonance positions for perturbed microresonator:

	TE ₀ resonances		
Direct simulation $n_c = 1.504$	1.021	1.044	1.068
Direct simulation $n_c = 1.5 + \text{Perturbational } \Delta\lambda$	1.021	1.044	1.068

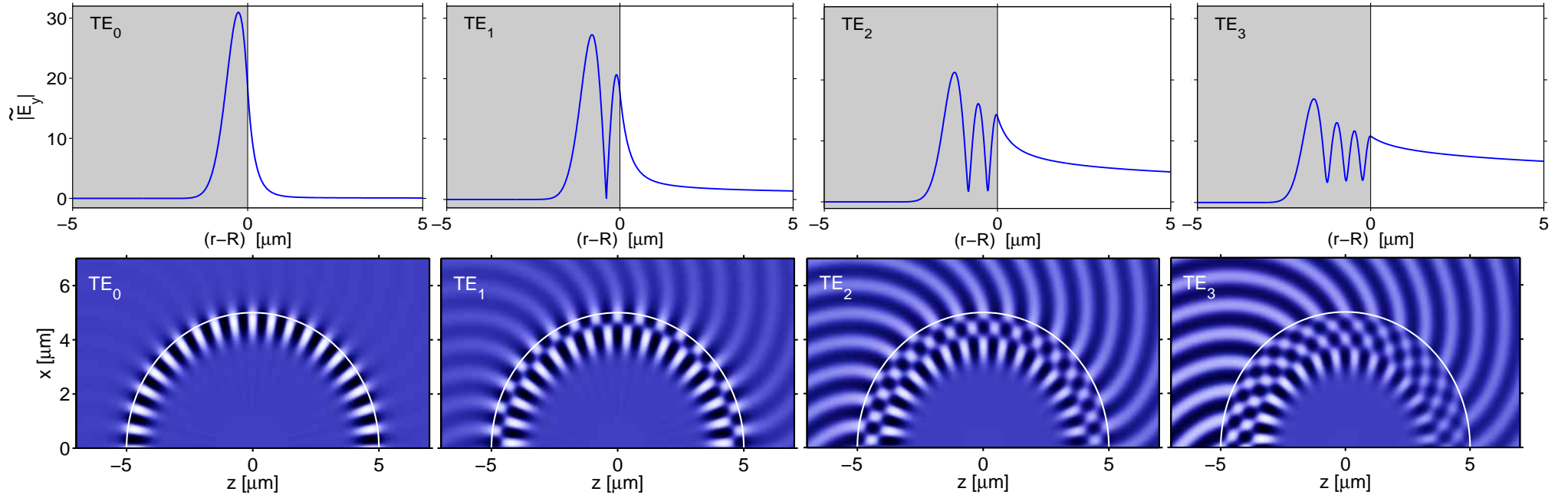
Whispering gallery modes



$|\tilde{E}_y|$ and real physical E_y field for a cavity with $n_c = 1.5$, $n_b = 1.0$, $R = 5 \mu\text{m}$, $\lambda = 1.05 \mu\text{m}$.

	$n_c = 1.5$
	N_{eff} (Direct)
TE ₀	$1.32793 - i 9.53121 \cdot 10^{-7}$
TE ₁	$1.16931 - i 4.03162 \cdot 10^{-4}$
TE ₂	$1.04222 - i 5.74103 \cdot 10^{-3}$
TE ₃	$0.92474 - i 1.31287 \cdot 10^{-2}$

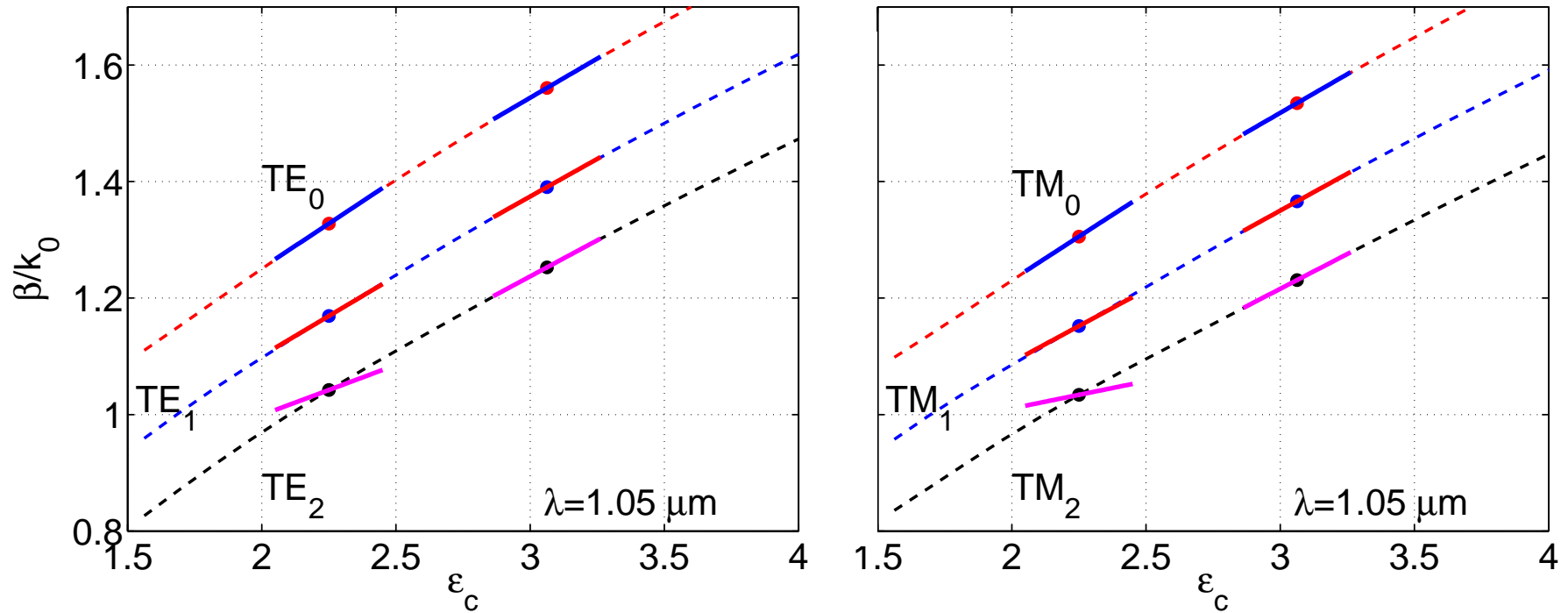
Whispering gallery modes



$|\tilde{E}_y|$ and real physical E_y field for a cavity with $n_c = 1.504$, $n_b = 1.0$, $R = 5 \mu\text{m}$, $\lambda = 1.05 \mu\text{m}$.

	$n_c = 1.5$	$n_c = 1.504$	
	N_{eff} (Direct)	N_{eff} (Direct)	$\Re(N_{\text{eff}})$ (Perturb.)
TE ₀	$1.32793 - i 9.53121 \cdot 10^{-7}$	$1.33160 - i 8.01245 \cdot 10^{-7}$	1.33161
TE ₁	$1.16931 - i 4.03162 \cdot 10^{-4}$	$1.17273 - i 3.61133 \cdot 10^{-4}$	1.17260
TE ₂	$1.04222 - i 5.74103 \cdot 10^{-3}$	$1.04555 - i 5.46820 \cdot 10^{-3}$	1.04429
TE ₃	$0.92474 - i 1.31287 \cdot 10^{-2}$	$0.92816 - i 1.28771 \cdot 10^{-2}$	0.92616

Bend mode phase shifts - II

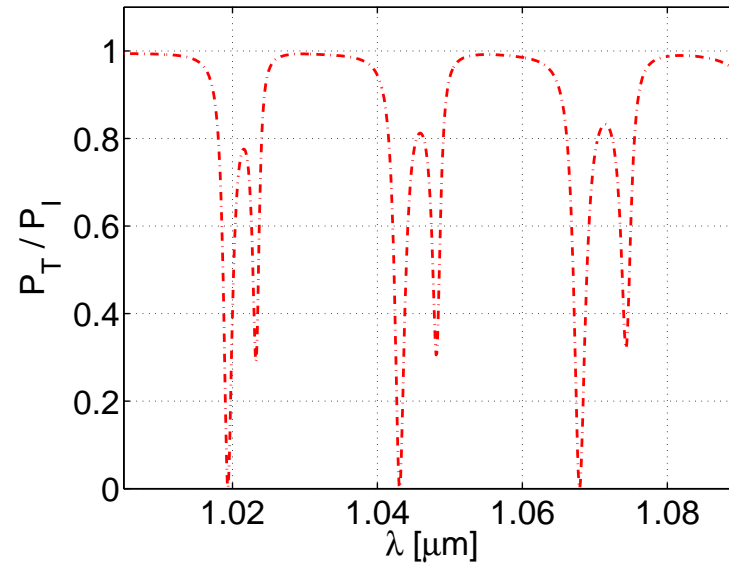


Dashed lines = β/k_0 direct calculations

Dots = Reference points ($\epsilon_c = (1.5)^2$ and $\epsilon_c = (1.75)^2$)

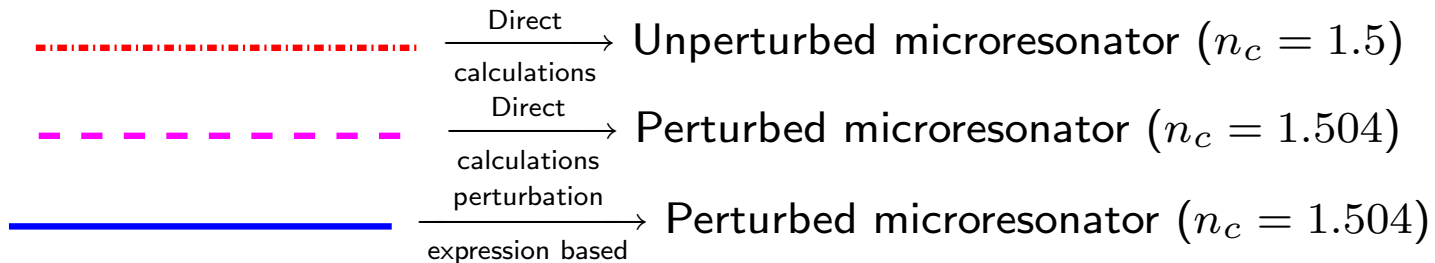
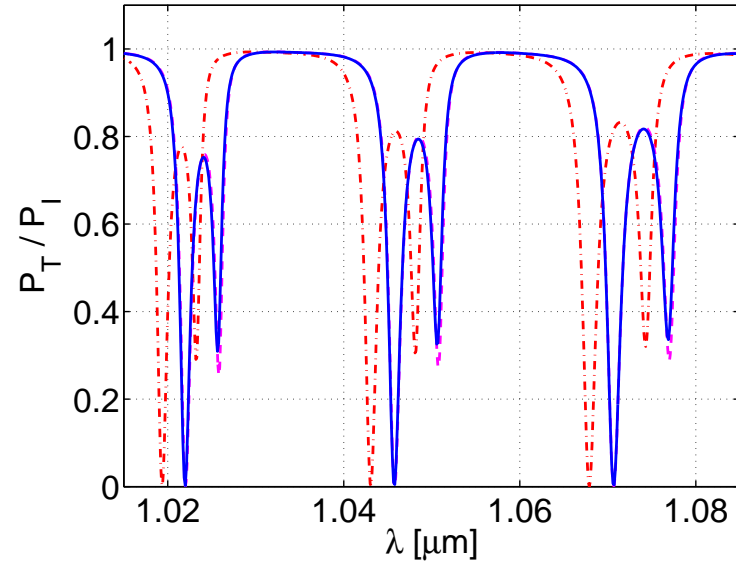
Solid segments = $\beta_{ref} + \delta\epsilon_c \left. \frac{\partial\beta}{\partial\epsilon_c} \right|_{ref}$, perturbational evolution of $\frac{\partial\beta}{\partial\epsilon_c}$

Microresonator spectrum perturbational evaluation - II



..... $\xrightarrow[\text{calculations}]{\text{Direct}}$ Unperturbed microresonator ($n_c = 1.5$)

Microresonator spectrum perturbational evaluation - II



- Resonance positions for perturbed microresonator:

	TE ₀ resonances			TE ₁ resonances		
	Direct simulation $n_c = 1.504$	1.022	1.026	1.046	1.051	1.071
Direct simulation $n_c = 1.5$ + Perturbational $\Delta\lambda$	1.022	1.026	1.046	1.051	1.071	1.077

Conclusions

- ✓ An expression for shifts in the bend mode phase constants due to localised perturbations of the permittivity distribution is derived.
- ✓ Using this expression, shifts of resonance wavelengths of microresonators due to perturbation of the cavity core refractive index are calculated.
- ✓ Shifts of the spectral response as predicted by the perturbational expression agree very well with the directly computed spectral response.



